

PARAMETER ANALYSIS OF A  
SINGLE STAGE INDUCTION MASS DRIVER

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ABSTRACT

The electromechanical behavior of a Single Stage Induction Mass Driver (SSIMD) has been modeled and analyzed with the aid of a computer. The analysis focused on the effects of circuit parameters on projectile velocity and system efficiency. A description of the model along with general rules for the design of efficient SSIMD systems are presented.

INTRODUCTION

The promise of achieving high velocities and high efficiencies has made the SSIMD an attractive concept for many applications. Since the mid-sixties the SSIMD has been adapted to such applications as: magnetic flux compressors [1], hypervelocity mass accelerators [2], induction reaction engines [3],[4], and electromagnetic metal formers [5].

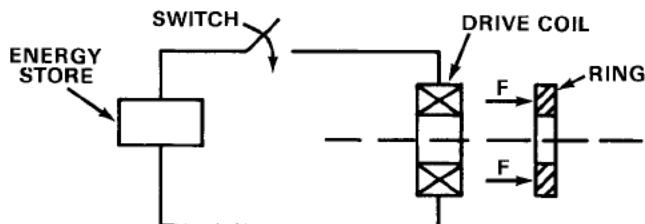


Figure 1: Single Stage Induction Mass Driver

The generic SSIMD (figure 1) basically consists of an energy storage (normally a capacitor bank), a switch, a drive coil, and a driven ring. The coil and ring are concentric and typically of equal diameter (the condition for maximum magnetic coupling). When a current pulse is applied to the drive-coil, a current is magnetically induced in the ring. The interaction between this ring current and the magnetic field (between the coil and the ring) produces a Lorentz Force which drives the ring.

The ring and coil geometries, and their relative orientation vary depending on the desired drive-direction (figure 2). In figure 2a the ring is driven axially, away from the drive coil. This is the geometry normally used in reaction engines and hypervelocity accelerators. In figure 2b the ring is driven inward radially, making this coil-ring arrangement favorable for flux-compression and metal forming applications.

In an effort to explore efficient induction driving, a twin capacitor bank system has been built and computer simulated. This twin bank scheme allows the use of high energy-density capacitors (for a compact system) by preventing voltage reversal across the banks. This paper is based mainly on experimental and simulation work with the above system.

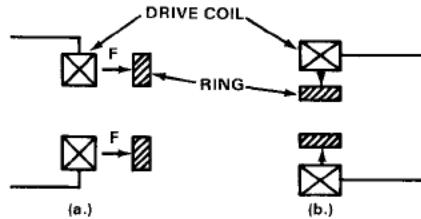


Figure 2: The two basic drive system geometries,  
a. typical geometry of axial accelerators  
b. typical geometry of flux compressors  
and metal formers

CIRCUIT MODEL

Our system has been electrically modeled by the circuit shown in figure 3. For discussion purposes, the circuit has been divided into two main sections: (1) the drive system, which includes the drive coil and the ring; and (2) the power system, which includes the capacitor banks, the switch and the power bus.

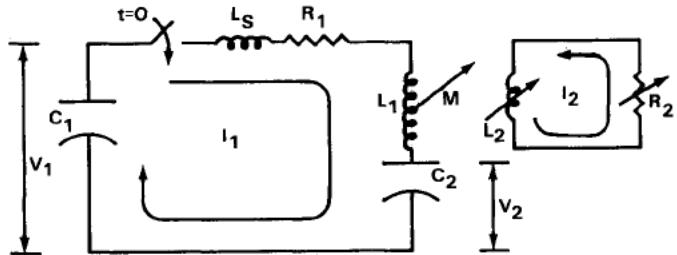


Figure 3: Circuit model of twin capacitor-bank device

The drive system has been modeled as a transformer with a shorted secondary and a mutual inductance (M) which is a function of ring separation. The self inductance of the primary ( $L_1$ ) is the drive coil inductance. This value is proportional to the coil mean diameter and the number of turns squared. The secondary self inductance ( $L_2$ ) is the ring inductance, which, within a few percent can be approximated by [6]

$$L_2 = \mu_0 R_m \left\{ \ln \left( \frac{R_m}{w} \right) + 0.9 \right\} \quad (1)$$

Where  $R_m$  is the ring's mean radius and  $w$  is its width.  $L_2$  is a variable quantity if ring deformation occurs during acceleration. Due to the typically high turns ratio between the primary and the secondary, the ring sees the highest currents in the system, and therefore experiences significant heating. For this reason,  $R_2$ , the ring resistance, is modeled as a temperature dependent element defined by the following equation:

$$R_2 = R_{20} \left\{ 1 + K_t (T - T_0) \right\} \quad (2)$$

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1. REPORT DATE <b>JUN 1985</b>	2. REPORT TYPE <b>N/A</b>	3. DATES COVERED <b>-</b>		
4. TITLE AND SUBTITLE  <b>Parameter Analysis Of A Single Stage Induction Mass Driver</b>			5a. CONTRACT NUMBER	
			5b. GRANT NUMBER	
			5c. PROGRAM ELEMENT NUMBER	
6. AUTHOR(S)			5d. PROJECT NUMBER	
			5e. TASK NUMBER	
			5f. WORK UNIT NUMBER	
7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES) <b>Naval Surface Weapons Center Dahlgren, Virginia 22448-5000</b>		8. PERFORMING ORGANIZATION REPORT NUMBER		
9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES)			10. SPONSOR/MONITOR'S ACRONYM(S)	
			11. SPONSOR/MONITOR'S REPORT NUMBER(S)	
12. DISTRIBUTION/AVAILABILITY STATEMENT <b>Approved for public release, distribution unlimited</b>				
13. SUPPLEMENTARY NOTES <b>See also ADM002371. 2013 IEEE Pulsed Power Conference, Digest of Technical Papers 1976-2013, and Abstracts of the 2013 IEEE International Conference on Plasma Science. Held in San Francisco, CA on 16-21 June 2013. U.S. Government or Federal Purpose Rights License.</b>				
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15. SUBJECT TERMS				
16. SECURITY CLASSIFICATION OF:			17. LIMITATION OF ABSTRACT <b>SAR</b>	18. NUMBER OF PAGES <b>4</b>
a REPORT <b>unclassified</b>	b ABSTRACT <b>unclassified</b>	c THIS PAGE <b>unclassified</b>	19a. NAME OF RESPONSIBLE PERSON	

Where  $K_t$  is the temperature coefficient of the ring-material (for copper and aluminum, at 20 °C,  $K_t = .0039$  [7]),  $T_0$  is the initial ring temperature,  $T$  is the operating temperature, and  $R_{20}$  is the initial ring resistance -defined by

$$R_{20} = \frac{2\pi R}{CwT_h} \quad (3)$$

where  $C$  is the ring-conductivity, and  $T_h$  is the ring thickness. The separation dependent mutual inductance,  $M$ , is determined using the following semiempirical procedure: First, the drive system circuit is transformed, ignoring  $R_2$ , to its T-model equivalent (figure 4).

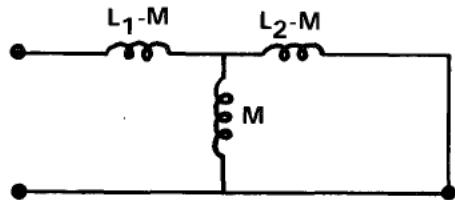


Figure 4: T-model of drive system used to obtain  $M(z)$

From this circuit, the effective drive system inductance,  $L_{eff}$ , (measured at the primary terminals) is

$$L_{eff} = L_1 - \left( \frac{M^2}{L_2} \right) \quad (4)$$

or

$$M = \sqrt{L_2 (L_1 - L_{eff})} \quad (5)$$

By measuring  $L_{eff}$  at increments of ring-coil normal separation,  $z$ , curves for  $L_{eff}$  and  $M$ , similar to the ones in figure 5, are obtained.

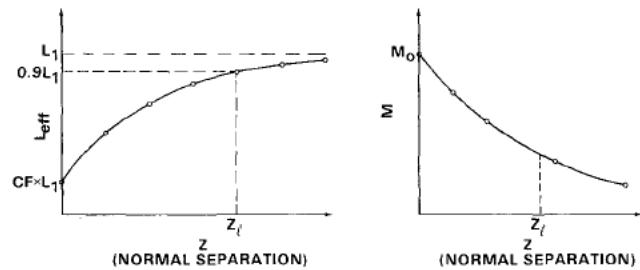


Figure 5: Experimental curves for drive system effective inductance and mutual inductance (coil mean diameter = 20.5 cm, 6 turns)

Finally, an exponential curve fit to the discrete  $M(z)$  and  $L_{eff}(z)$  yields the following expressions

$$L_{eff} = L_1 \left\{ 1 - \left( \frac{L_{20}}{L_1} \right) K_0^2 e^{-2z/z_\ell} \right\} \quad (6)$$

$$M = \sqrt{\frac{L_{20} L_1}{L_1}} K_0 e^{-z/z_\ell} \quad (7)$$

where  $K_0$  is the initial (at  $z=0$ ) coupling coefficient (as used in transformer terminology) and is defined by

$$K_0 = \sqrt{1 - C_f} \quad (8)$$

where  $C_f$ , the compensation factor, is the factor by which  $L_1$  is reduced when the ring is at  $z=0$ .  $z_\ell$ , the coupling length, is the separation up to which

there is significant coil-ring magnetic interaction, and it is defined as the separation when  $L_{eff} = 0.9L_1$ . This coupling length has experimentally been found to roughly equal .17D, where D is the mean drive-coil diameter. This result agrees with streak camera observations by Bondaletoff [2] of the maximum separation up to which the ring is accelerated.

In the power-system,  $R_1$  is the lumped representation of the primary loop resistance, and it is obtained from the exponential decay of the oscillating primary current ( $I_1$ ) without the ring in place. Similarly, the stray inductance,  $L_s$ , is obtained from the oscillating frequency of  $I_1$  when the coil is replaced by a shorting strap. Finally,  $C_1$  and  $C_2$  represent two capacitor banks of equal capacitance; where initially  $C_1$  is charged and  $C_2$  is uncharged.

### SIMULATION

In order to predict the behavior of our SSIMD, a computer simulation has been developed using the circuit model previously described. The time behavior of the following variables were of interest: the primary current ( $I_1$ ), the secondary (ring) current ( $I_2$ ), the capacitor banks voltages ( $V_1$  and  $V_2$ ), the ring's normal velocity ( $v$ ), the normal force exerted on the ring ( $F$ ), the normal ring position ( $z$ ), and the heat energy dissipated in the ring ( $E_2$ ). The approach taken was to first write differential loop equations for the primary and the secondary of figure 3, and then manipulate them into a form that could be solved numerically. For the primary

$$L_1 \dot{I}_1 + I_1 R_1 + Q_2 / C_2 - Q_1 / C_1 + M \dot{I}_2 + I_2 \dot{M} = 0 \quad (9)$$

and the secondary

$$(L_2 + R_2) \dot{I}_2 + L_2 \dot{I}_1 + M \dot{I}_1 + I_1 \dot{M} = 0 \quad (10)$$

Where  $Q_1$  and  $Q_2$  are respectively the charge stored in  $C_1$  and  $C_2$ . The numerical technique used was a fourth order Runge-Kuta Method, which requires that the derivative of the variable to be solved be equal to an expression which does not include that derivative. By respectively multiplying equations (9) and (10) by  $L_2$  and  $M$ , then subtracting the resultant equations an expression for  $dI_1/dt$  is obtained

$$\dot{I}_1 = \frac{1}{L_2 L_1 - M} \{ Q_1 L_2 / C_1 - L_2 Q_2 / C_2 + (M \dot{M} - L_2 R_2) I_1 + (M \dot{L}_2 + R_2 M - L_2 \dot{M}) I_2 \} \quad (11)$$

Similarly, by respectively multiplying equations (9) and (10) by  $M$  and  $L_1$ , and then subtracting the resultant equations  $dI_2/dt$  is obtained

$$\dot{I}_2 = \frac{1}{L_2 L_1 - M} \{ M Q_2 / C_2 - M Q_1 / C_1 + (M R_1 - M \dot{L}_1) I_1 + (M \dot{M} - L_1 \dot{L}_2 - R_2 L_1) I_2 \} \quad (12)$$

The rest of the necessary equations are:

$$\dot{Q}_1 = -I_1 \quad (13)$$

where

$$V_1 = V_0 + Q_1 / C_1 \quad (14)$$

$$Q_2 = I_1$$

(15)

where

$$V_2 = Q_2 / C_2$$

(16)

The force exerted on the ring is defined by [8]

$$F = I_1 I_2 dM/dz$$

or

$$\dot{V} = F/m = (I_1 I_2 / m) (dM/dz)$$

(18)

where  $m$  is the mass of the ring.

$$\dot{z} = V$$

and

$$\dot{E}_2 = I_2^2 R_2$$

Figure 6 shows the various simulation output curves for a sample run along with comparable measured variables.

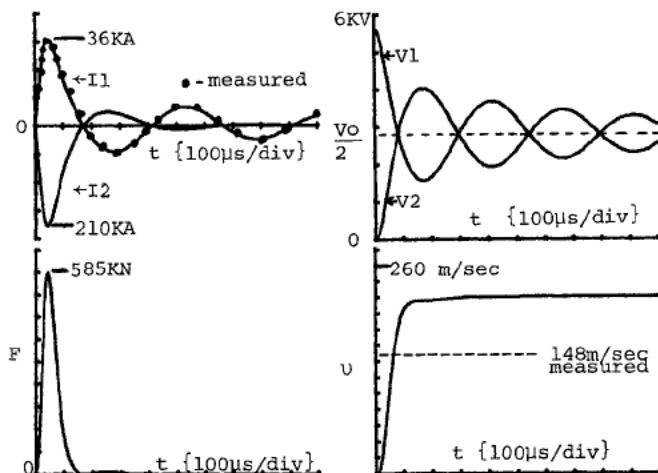


Figure 6: Computer generated curves compare to measurements for the following case:  
 $L_1 = 14\mu H$ ,  $C_1 = C_2 = 828\mu F$ ,  $R_1 = 34m\Omega$ ,  $L_s = 1\mu H$ ,  
 $Z_g = .035m$ ,  $L_{2o} = .32\mu H$ ,  $R_{2o} = .65m\Omega$ ,  $m = .115$  Kg,  
and  $K_0 = .89$

In general, good agreement exists between measured and calculated values. The most significant error exists in the velocity prediction (typically by a factor of  $\sim 0.7$ ). This error is probably due to aerodynamic effects which have not been modeled.

#### ANALYSIS

In an effort to empirically develop rules for the design of energy efficient SSIMD, the previously described simulation was used to generate the family of curves shown in figure 7.

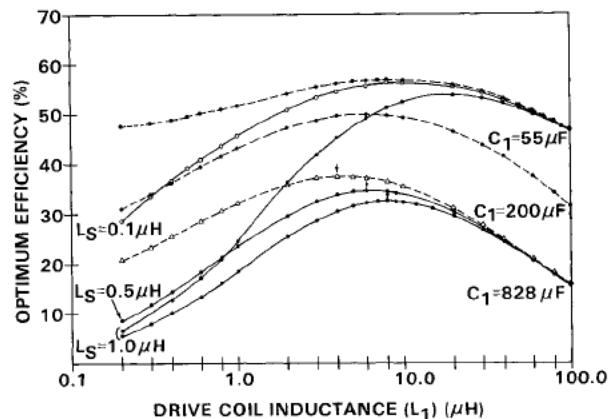


Figure 7: Computer generated optimum efficiency curves as a function of  $L_1$ ,  $C_{eff}$ , and  $L_s$ .  $R_1 = 15.6m\Omega$ ,  $R_2 = 1.4m\Omega$ ,  $L_2 = .32\mu H$ ,  $K_0 = .89$ ,  $Z_g = .035m$ , and  $m = .1$  Kg

The solid and dashed curves respectively represent the presence and absence of stray inductance in the system, and have been generated for three values of  $C_1$  ( $C_2 = C_1$ ).

Each point in these curves represents an optimum efficiency for a given set of circuit elements, and was found by stepping through the charge voltage until an efficiency maximum was reached. The purpose of changing the voltage was to vary the ring velocity. These points agree with the following criterion [4]:

$$\langle v \rangle = z_g / \langle T/2 \rangle \quad (21)$$

where  $\langle v \rangle$  is the average ring velocity and  $\langle T/2 \rangle$  is the average first half-period of the primary current - defined by

$$\langle T/2 \rangle = \pi \sqrt{L_{eff} C_{eff}} \quad (22)$$

where

$$\langle L_{eff} \rangle = \frac{1}{z_g} \int_0^{z_g} L_{eff} dz + L_s \quad (23)$$

and  $C_{eff}$  is the effective system capacitance. For example, in the twin capacitor bank system  $C_{eff}$  equals the series combination of  $C_1$  and  $C_2$ . The above criterion simply states that in order to drive a ring most efficiently, the ring must remain in the coupling region (between  $z=0$  and  $z=z_g$ ) for exactly  $\langle T/2 \rangle$ .

By comparing the solid and dashed curves the adverse effects of stray inductance on efficiency is clearly seen. When  $L_s$  is of the order of  $L_1$  a significant energy division occurs between the two: this division reduces the amount of energy available at the drive coil, and consequently reduces drive efficiency. Aside from the effects of  $L_s$  on efficiency, when  $L_1$  is made too small, the efficiency is also reduced, as shown by the dashed curves of figure 7. This effect can be attributed to the effective coupling parameter  $dM/dz|_{eff}$ , which is present in the force equation, as shown below

$$F = I_1^2 I_2^2 \frac{dM}{dz} = I_1^2 M/L_2 \frac{dM}{dz} = I_1^2 \frac{dM}{dz} \Big|_{\text{eff}} \quad (24)$$

This effective coupling parameter is used as a figure of merit when comparing different induction accelerators [4]. The dependence of  $dM/dz|_{\text{eff}}$  on  $L_1$  can be seen when  $M$  and  $dM/dz$  are spatially averaged from,  $z=0$  to  $z_L$  and substituted into eqn. (24)

$$\left\langle \frac{dM}{dz} \right\rangle_{\text{eff}} = \frac{.4K_o^2 L_1}{z_L} \quad (25)$$

Although it is undesirable for  $L_1$  to be small, it is also undesirable for it to be too large, since  $I_1$ , (the other term in eqn. (24)), will be reduced.

These curves also show the beneficial effects, in terms of efficiency, of reducing the effective system capacitance. Since lower capacitance translates to higher velocity (from the optimum efficiency criteria), the trends in figure 7 imply that induction drivers have a natural tendency to efficiently drive rings to high velocities. Unfortunately, the price paid for reducing the capacitance and still operate at optimum efficiency is an increase in charge voltage, which ultimately becomes one of the practical limits on device performance. This voltage limit is approximately 100Kv [4]. An estimate of the optimum charge voltage, which agrees with the simulation results, can be derived from the force equation (eqn. (17)) by averaging all of the pertinent variables: if  $L_2/R_2 \gg \langle T/2 \rangle$ , which usually is the case, eqn. (17) can be rewritten as

$$F = I_1^2 (M/L_2) \frac{dM}{dz} = ma \quad (26)$$

Averaging this expression and replacing  $\langle a \rangle$ , the average ring acceleration, with  $\langle v \rangle / \langle T/2 \rangle$

$$(2m \langle v \rangle) / \langle T/2 \rangle = I_p^2 / 4 \langle M/L_2 \frac{dM}{dz} \rangle \quad (27)$$

where  $I_1$  was replaced by its RMS value from  $t=0$  to  $t=\langle T/2 \rangle$ ,  $I_p/\sqrt{4}$ .  $I_p$  is the peak primary current, defined by

$$I_p = V_o \sqrt{\frac{C_{\text{eff}}}{\langle L_{\text{eff}} \rangle}} \quad (28)$$

Finally, combining eqns. (25) and (26) yields an expression for the optimum charge voltage.

$$V_o = \sqrt{\frac{4m \langle v \rangle}{\pi}} \frac{\sqrt{\langle L_{\text{eff}} \rangle}}{\left( C_{\text{eff}} \right)^{1.5} \left\langle \frac{M}{L_2} \frac{dM}{dz} \right\rangle} \quad (29)$$

The curves in figure 7 are independent of ring mass, but the charge voltage required to operate at optimum efficiency is proportional to the square root of the mass -as indicated by Eqn. (29).

#### DESIGN PROCEDURE

The purpose of this procedure is to determine the values of the circuit elements and the charge voltage, based on a design velocity and drive system

dimensions, which will yield an efficient SSIMD. The objective is to choose the circuit parameters which will cause the device to operate near the optimum efficiency peaks of figure 7. The procedure assumes a capacitive discharge, series RLC-circuit, and a drive-coil and ring of roughly equal diameter. The first step is to minimize the values of  $L_s$  and  $R_1$ .  $L_s$  is determined by the geometry of the power system; some typical low values for parallel plate and coaxial geometries are respectively 250nH[3] and 120nH 4 (at voltages near 30Kv).  $R_1$  is the conductor resistance of the primary circuit, and is typically of the order of 10E-3 ohms. The application normally dictates the dimensions of the drive-coil and the ring. These dimensions in turn determine  $L_2$  and the lower limit of  $L_1$ . By increasing the number of turns,  $L_1$  is increased until  $L_1 > 10L_s$ . The upper limit on  $L_1$  is not easily defined, and it may be necessary to use a computer model to determine its optimum value.  $R_2$  is determined by the ring's dimensions and material properties. This parameter has a strong impact on efficiency through the  $I_2^2 R_2$  losses, and therefore should be minimized.  $C_{\text{eff}}$  is chosen according to the optimum efficiency criteria (eqn. (21)), where the average velocity can be approximated by one half the design velocity. Finally, a rough estimate of the required voltage is obtained using eqn. (29), where the spatial averages for  $M$ ,  $L_{\text{eff}}$ , and  $dM/dz$  are found from their previously described equations.

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